

Instantons defined by Lie groups

Simon Salamon

12 August 2014

Hannover

M.F. Atiyah, D.G. Drinfeld, N.J. Hitchin, Y.I. Manin: *Construction of instantons*, Phys. Lett. 65A (1978), 185–7

R.S. Ward: *Completely solvable gauge field equations in dimension greater than four*, Nucl. Phys. B236 (1984), 381–96

M. Mamone Capria, S.M. Salamon: *Yang-Mills fields on quaternionic spaces*, Nonlinearity 1 (1988) 517–530

R. Reyes Carrión: *A generalization of the notion of instanton*
Differ. Geom. Appl. 8 (1998), 1–20

4-dimensional origins

On an oriented Riemannian 4-manifold,

$$\Lambda^2 T^* M = \Lambda_+^2 \oplus \Lambda_-^2$$

since $\mathfrak{so}(4) = \mathfrak{su}(2)_+ \oplus \mathfrak{su}(2)_-$, and there is an elliptic complex

$$0 \rightarrow \Omega^0 \xrightarrow{d} \Omega^1 \xrightarrow{d_-} \Omega_-^2 \rightarrow 0.$$

If Q is a principal $SU(2)$ -bundle with self-dual connection A with (so $F = dA + A \wedge A \in \Omega_+^2$ and $*F = F$) then H^1 of the complex

$$0 \rightarrow \Omega^0(\text{ad } Q) \rightarrow \Omega^1(\text{ad } Q) \rightarrow \Omega_-^2(\text{ad } Q) \rightarrow 0$$

captures infinitesimal deformations modulo gauge equivalence.

The index is $8k - 3$, and a framed moduli space of dimension $8k$ is hyperkähler (meaning holonomy $Sp(k)$), given by HK quotients $\mathbb{H}^{k(k+3)/2} // O(k) \cong \mu_\infty^{-1}(0) / \mathcal{G}$ [D].

More generally, for $G \subset \mathrm{SO}(n)$, we can write

$$\Lambda^2 \cong \mathfrak{so}(n) = \mathfrak{g} \oplus \mathfrak{g}^\perp.$$

Given a G -structure on M (in fact, an $N(G)$ -structure) this decomposition passes to 2-forms.

Definition. In this context, an *instanton* is a connection (on a bundle over M) whose curvature 2-forms F_i^j lie in \mathfrak{g} .

Example. If the holonomy reduces to G then the Levi-Civita connection is an instanton because R_{ijkl} belongs to the kernel of

$$S^2(\mathfrak{g}) \subset \mathfrak{g} \otimes \Lambda^2 \subset \Lambda^2 \otimes \Lambda^2 \longrightarrow \Lambda^4.$$

On the other hand, the Killing form in $S^2(\mathfrak{g})$ maps to a *non-zero* 4-form unless it represents curvature of a Riemannian symmetric space. So 4-forms arise in abundance!

Hitchin-Kobayashi correspondence

$\frac{4}{19}$

If $G = \mathrm{SU}(n) \subset \mathrm{SO}(2n)$ so $N(G) = \mathrm{U}(n)$ then

$$\Lambda^2 = [[\Lambda^{2,0}]] \oplus [\Lambda_0^{1,1}] \oplus \langle \omega \rangle,$$

and $\mathfrak{g} = \mathfrak{su}(n) \cong [\Lambda_0^{1,1}]$. An instanton is a connection with $(1,1)$ curvature and vanishing trace $F \wedge \omega^{n-1}$, though this would force $c_1 = 0$. More generally we require that the trace be a (constant) multiple of the identity, the Hermitian-Einstein condition.

Over a complex manifold:

- A connection with $(1,1)$ curvature on a vector bundle renders it a *holomorphic* bundle [cf. NN].
- A holomorphic bundle admits a unique connection with $(1,1)$ curvature compatible with a given *Hermitian metric* on its fibres.

Theorem [D,UY]. On a compact Kähler manifold, an irreducible holomorphic vector bundle admits a HE connection iff it is *stable*.

Using a 3-form

If $G = G_2 \subset SO(7)$ so $N(G) = G$ then

$$\begin{aligned}\Lambda^2 &= \Lambda_{14}^2 \oplus \Lambda_7^2 && \cong \mathfrak{g}_2 \oplus \Lambda^1 \\ \Lambda^3 &= \Lambda_{27}^3 \oplus \Lambda_7^3 \oplus \langle \varphi \rangle && \cong S_0^2 \oplus \Lambda^1 \oplus \mathbb{R}.\end{aligned}$$

Example. If $\varphi = (12 - 34)5 + (13 - 42)6 + (14 - 23)7 + 567$, we have $12+34, 13+24, 14+23 \in \Lambda_+^2 \subset \mathfrak{g}_2$.

An instanton is characterized by the equivalent equations

$$F_7 = 0, \quad F \wedge (*\varphi) = 0, \quad F \wedge \varphi = *F.$$

Instantons are YM connections because

$$c_2(F) \cup [\varphi] = \int F \wedge F \wedge \varphi = 4\|F_{14}\|^2 - 18\|F_7\|^2$$

and $\|F\|^2$ has an absolute minimum if $F_7 = 0$.

grad, curl, div in dim 7

Suppose that M^7 has a G_2 structure. Consider

$$0 \rightarrow \Omega^0 \xrightarrow{d} \Omega^1 \xrightarrow{D_1} \Omega^2_7 \xrightarrow{D_2} \Omega^3_1 \rightarrow 0$$

where $D_1 = \pi_7 \circ d$ arises from the cross product. It is complex iff

$$D_2 \circ D_1 = 0 \iff d(\Omega^2_{14}) \subseteq \langle \varphi \rangle^\perp \iff d * \varphi = 0.$$

Lemma [CN]. If M^7 is oriented and spin it has a G_2 structure, indeed one with $d * \varphi = 0$.

Given an instanton on a bundle Q , we can extend the operators so $D \circ D = F_7 = 0$ and obtain an elliptic complex

$$0 \rightarrow \Omega^0(\text{ad } Q) \rightarrow \Omega^1(\text{ad } Q) \rightarrow \Omega^1(\text{ad } Q) \rightarrow \Omega^0(\text{ad } Q) \rightarrow 0$$

whose H^1 parametrizes infinitesimal deformations. Close analogue with the de Rham complex $1 \rightarrow 3 \rightarrow 3 \rightarrow 1$ over a 3-manifold.

Integrability

$\frac{7}{19}$

We shall construct a differential complex for any $G \subset \mathrm{SO}(n)$ that begins $0 \rightarrow \Omega^0 \rightarrow \Omega^1 \rightarrow \dots$

Set $\Lambda^{-2} = \Lambda^{-1} = 0$ and $A^k = (\mathfrak{g} \wedge \Lambda^{k-2})^\perp$. Define

$$D: \mathcal{A}^k \subset \Omega^k \xrightarrow{d} \Omega^{k+1} \xrightarrow{\pi} \mathcal{A}^{k+1}.$$

Here $\mathcal{A}^k = \Gamma(M, A^k)$, so $\mathcal{A}^k = \Omega^k$ for $k = 0, 1$.

Proposition. $D^2 = 0$ if and only if $d: \Omega^2 \rightarrow \Omega^3$ maps sections of \mathfrak{g} to those of $\mathfrak{g} \wedge \Lambda^1$.

This is obviously true if the holonomy of M lies in $N(G)$; in general it is a condition on the *intrinsic torsion* $\tau \in \Gamma(M, \mathfrak{g} \otimes \Lambda^1)$.

In any case, we would like

$$0 \rightarrow \mathcal{A}^0 \rightarrow \mathcal{A}^1 \rightarrow \mathcal{A}^2 \rightarrow \dots \rightarrow 0$$

to be *elliptic*. It is when G equals $\mathrm{SU}(n), \mathrm{G}_2, \mathrm{Spin} 7, \mathrm{Sp}(n), \dots$

Nearly-Kähler 6-manifolds

If $G = \mathrm{SU}(3)$ so $N(G) = \mathrm{U}(3)$, the complex becomes

$$0 \rightarrow \Omega^0 \rightarrow \Omega^1 \rightarrow \Gamma(\llbracket \Lambda^{2,0} \rrbracket \oplus \langle \omega \rangle) \rightarrow \Gamma(\llbracket \Lambda^{3,0} \rrbracket) \rightarrow 0.$$

with dimensions $1 \rightarrow 6 \rightarrow 7 \rightarrow 2$, provided *most* of the Nijenhuis tensor vanishes! We get a theory of instantons over nearly-Kähler 6-manifolds (meaning $(\nabla_X J)X = 0$).

Example. The twistor spaces $\mathbb{C}\mathbb{P}^3 \rightarrow S^4$ and $\mathbb{F}^3 \rightarrow \mathbb{C}\mathbb{P}^2$ both have a $\mathrm{U}(1)$ -connection A_1 whose curvature F_1 is a Kähler form and another NK 2-form F_2 such that $F_2 \wedge F_2 = d\psi$ with $F_1 \wedge \psi = 0$.

$$\implies 0 = d(F_1 \wedge \psi) = dF_1 \wedge \psi + F_1 \wedge (F_2 \wedge F_2) = 2F_1 \wedge (*F_2).$$

Thus A_1 is an instanton for the NK metric.

If $G = \mathrm{Sp}(n) \subset \mathrm{SO}(4n)$ so $N(G) = \mathrm{Sp}(n)\mathrm{Sp}(1)$ and

$$\begin{aligned}(\Lambda^1)_c &= E \otimes H \quad (\text{cf. } S \otimes \tilde{S}) \\(\Lambda^2)_c &= S^2 E \oplus S^2 H \oplus (\Lambda_0^2 E \otimes S^2 H) \\ &\cong \mathfrak{sp}(n) \oplus \mathfrak{sp}(1) \oplus \mathfrak{m}.\end{aligned}$$

Manifolds M^{4n} ($n \geq 2$) like $\mathbb{H}\mathbb{P}^n$ with holonomy in $N(G)$ are quaternion-Kähler and behave as if they were *nearly hyperkähler*.

Since $S^2 E = \Lambda_I^{1,1} \cap \Lambda_J^{1,1} \oplus \Lambda_K^{1,1}$, instantons give rise to holomorphic bundles over the twistor space Z^{2n+1} , which fibres over M .

Using a 4-form again, one shows that the Yang-Mills functional has a critical point whenever any 2 of the 3 components vanish.

Not an abs max/min if $F \in \Gamma(M, \mathfrak{m})$, but no examples known.

Instantons via quaternions

10
19

Take $M^{4n} = \mathbb{H}\mathbb{P}^n$. Let $\mathbf{q} = (q_0, q_1, \dots, q_n) \in \mathbb{H}^{n+1} \setminus 0$, $m = [\mathbf{q}]$.

A linear form $\sum a_r q_r$ (with $a_r \in \mathbb{H}$) defines a section of the tautological line bundle H (fibre $\mathbb{H} = \mathbb{C}^2$) inside $\underline{\mathbb{H}}^{n+1}$, and

$$E_m = \ker(\mathbf{q}^\top : \mathbb{H}^{n+1} \rightarrow H_m), \quad \text{so } E = H^\perp.$$

Take matrices $A_0, A_1, \dots, A_n \in \mathbb{H}^{n+k, k}$ and set $\mathbb{A} = \sum_{r=0}^n A_r q_r$.

Theorem If $A_r^* A_s$ is symmetric for all r, s and \mathbb{A} has rank k for all $\mathbf{q} \neq 0$ then $\ker \mathbb{A}$ is an $\mathrm{Sp}(n)$ instanton on $\mathbb{H}\mathbb{P}^n$.

Proof. Relies on the fact that the real components of

$$dq_r \wedge d\bar{q}_s = (du_r + j dv_r) \wedge (d\bar{u}_r - d\bar{v}_r j)$$

span the subspace $\mathfrak{sp}(n)$ of Λ^2 . Projection to $\ker \mathbb{A}$ equals $\Pi = 1 - \mathbb{A}(\mathbb{A}^* \mathbb{A})^{-1} \mathbb{A}^*$ and the induced curvature is $\Pi(d\Pi \wedge d\Pi)\Pi$.

The twistor space

11
19

... of $\mathbb{H}\mathbb{P}^n$ is $\mathbb{C}\mathbb{P}^{2n+1}$, which is the total space of

$$\begin{array}{c} \mathbb{P}_c(H) \\ \downarrow \\ \mathbb{H}\mathbb{P}^n. \end{array}$$

The instantons $F = \ker \mathbb{A}$ pull back to holomorphic bundles $\pi^* F$ (fibre \mathbb{C}^{2n}) with $c(F) = (1-h)^{-k}$, characterized by

$$H^q(\mathbb{C}\mathbb{P}^{2n+1}, \pi^* F \otimes \mathcal{O}(p)) = 0 \quad \begin{cases} q = 1, & p \leq -2 \\ 2 \leq q \leq n, & p \in \mathbb{Z}. \end{cases}$$

Example. For $n = k = 2$, we can take $\mathbb{A} = \begin{pmatrix} q_0 & 0 \\ 0 & q_0 + q_2 \\ q_1 & q_2 \\ q_2 & q_1 \end{pmatrix}$.

The deformation complex for $n = 1$ has $h^1 = 8k - 3$.

Proposition. If $n = 2$ the deformation complex of the instantons above has $h^1 - h^2 = \frac{3}{2}k(17-k) - 10 = 14, 35, 53, \dots$

4-forms in 8 dims

Many interesting geometries in dim 8 are characterized by 4-forms: elements of $\Lambda^4(\mathbb{R}^8)^*$, the isotropy representation of the symmetric space $E^7/SU(8)$. The complicated orbit structure for $SL(8, \mathbb{R})$ acting on the 4-forms can be understood via roots of E_7 [M].

We focus on the inclusions

$$\begin{array}{ccccc} & \nearrow & \text{Sp}(2)\text{Sp}(1) & \searrow & \\ \text{Sp}(2) & & & & \text{SO}(8). \\ & \searrow & \text{Spin } 7 & \nearrow & \end{array}$$

$\text{Sp}(2)$ fixes a triple $\omega_1, \omega_2, \omega_3$, whilst

$$\begin{array}{ll} \text{Sp}(2)\text{Sp}(1) & \text{stabilizes } \Omega = \omega_1 \wedge \omega_1 + \omega_2 \wedge \omega_2 + \omega_3 \wedge \omega_3 \\ \text{Spin } 7 & \text{stabilizes } \Phi = -\omega_1 \wedge \omega_1 + \omega_2 \wedge \omega_2 + \omega_3 \wedge \omega_3. \end{array}$$

We shall investigate the topology defined by the two rank 3 groups.

Proposition. If a compact, oriented M^8 has a Spin 7 or a $\mathrm{Sp}(2)\mathrm{Sp}(1)$ structure then M is spin and $8\chi = 4p_2 - p_1^2$.

Proof.

$$\Delta_+ \cong \begin{cases} \Lambda^0 \oplus \Lambda_7^2 & = 1 + 7 & \text{for Spin 7} \\ S^2H \oplus \Lambda_0^2 E & = 3 + 5 & \text{for Sp(2)Sp(1)} \end{cases}$$

and in both cases $\Delta_- \cong \Lambda^1 \cong TM$. The Euler class e satisfies

$$\mathrm{ch}(\Delta_+ - \Delta_-) = e \hat{A}^{-1} = e \left(1 - \frac{1}{24}p_1 + \hat{A}_2 + \dots\right)^{-1}$$

where $\hat{A}_2 = \frac{1}{5760}(7p_1^2 - 4p_2)$.

Theorem. If $\mathrm{Hol}(M) \subseteq \mathrm{Sp}(2)\mathrm{Sp}(1)$ and $s > 0$ then $\hat{A}_2 = 0$.
If $\mathrm{Hol}(M) = \mathrm{Spin} 7$ then $\hat{A}_2 = 1$.

So the QK 8-manifolds $\mathbb{H}\mathbb{P}^2$, $G_2/SO(4)$, $\mathbb{G}r_2(\mathbb{C}^4)$ all admit a Spin 7 structure but not the rival holonomy!

The remarkable space $G_2/SO(4)$

$\frac{14}{19}$

- parametrizes coassociative 4-planes in $\mathbb{R}^7 \subset \mathbb{O}$ [HL].
- As an application, its orbits under $SO(4)$ are relevant to the classification of coassociative submanifolds of the G_2 manifold $\Lambda_-^2(S^4)$ that are deformations of T^*S^2 [KS].
- Removing $\mathbb{C}P^2$ and quotienting out by \mathbb{Z}_3 , we get \mathcal{N}/\mathbb{H}^* , where \mathcal{N} is the principal nilpotent orbit in $\mathfrak{sl}(3, \mathbb{C})$. The latter is HK and the quotient QK [K,S].
- There is a construction of QK metrics from 5-manifolds with generic 2-plane distributions that are modelled asymptotically on the noncompact dual $G_2^s/SO(4)$ [L,B,D].

A Dirac complex

If M^8 has a Spin 7 structure then (\mathcal{A}^\bullet, D) coincides with

$$0 \rightarrow \Omega^0 \rightarrow \Omega^1 \rightarrow \Omega_7^2 \rightarrow 0.$$

Only if it is written backwards do we need the holonomy condition!

It is easy to compute the deformation index

$$h^0 - h^1 + h^2 = \int \text{ch}(\text{ad } Q) \hat{A}$$

though h^2 is again unknown.

Example. For the Levi-Civita instanton Λ^1 on a manifold with holonomy equal to Spin 7, it equals $8 - \frac{1}{3}\chi$. The latter is an integer because

$$25 + b_2 + 2b_4^- = b_3 + b_4^+.$$

... of a $\mathrm{Sp}(2)\mathrm{Sp}(1)$ structure lies in $\Lambda^1 \otimes \mathfrak{m} =$

EH	KH
ES^3H	KS^3H

M^8 is *quaternionic* if $\tau \in$ row 1; it is *ideal* if $\tau \in$ col 1.

Surprise. $\mathrm{SU}(3)$ possesses invariant structures of both types [J,M].

If M^8 is quaternionic, there is an elliptic complex

$$0 \rightarrow \Gamma(S^2H) \rightarrow \Gamma(EH) \rightarrow \Gamma(\Lambda_0^2E) \rightarrow 0,$$

which corresponds to the sheaf $\mathcal{O}(-2)$ on the twistor space Z^5 .

Tensor by S^2H to obtain (\mathcal{A}^\bullet, D) . Passing to $\mathcal{O}(-3)$ gives

$$0 \rightarrow \Gamma(H) \xrightarrow{\partial} \Gamma(E) \xrightarrow{\square} \Gamma(E) \rightarrow \Gamma(H) \rightarrow 0$$

where ∂ is a Fueter operator and \square is second order [B].

Some K theory

Associated to the Dirac operator on $\mathbb{H}\mathbb{P}^2$ is the virtual vector bundle

$$\Lambda_0^2(E - H) = \Lambda_0^2 E - E \otimes H + S^2 H.$$

Now, $E - H$ can't be a genuine vector bundle because:

- any monomorphism $H \rightarrow E$ would define a nowhere zero section of $E \otimes H \cong T\mathbb{H}\mathbb{P}^2$ but $\chi = 3$;
- $E - H$ has rank 2, but a calculation shows $c_4(E - H) \neq 0$!

In fact, $c(H) = 1 - h$ so

$$c(E - H) = c(\underline{\mathbb{C}}^6 - 2H) = (1 - h)^{-2} = 1 + 2h + 3h^2.$$

By contrast,

$$c(\Lambda_0^2 E - H) = c(\Delta_+ - S^2 H - H) = 1 - 3h.$$

We shall see that this time the difference is a vector bundle.

Horrocks's instanton

$\frac{18}{19}$

Theorem. There exists a rank 3 complex vector bundle V over $\mathbb{H}P^2$ with $c_2=3h$, and an $SU(3)$ -connection with $F \in \Gamma(\mathfrak{sp}(2))$.

Proof. Recall that $E = \ker(p_1: \underline{\mathbb{C}}^6 \rightarrow H)$. Similarly,

$$\Lambda_0^2 E = \ker(\Lambda_0^2(\underline{\mathbb{C}}^6) \rightarrow \underline{\mathbb{C}}^6 \wedge H).$$

Fix a reduction of $Sp(3)$ to $SU(3)$, giving $\underline{\mathbb{C}}^6 = \Lambda^{1,0} \oplus \Lambda^{0,1}$ and

$$p_2: \Lambda_0^2(\underline{\mathbb{C}}^6) \rightarrow \underline{\mathbb{C}}^6.$$

Then $p_1 \circ p_2$ has rank 2 everywhere. The instanton connection on $V = \ker(p_1 \circ p_2)$ is induced from that on $\Lambda_0^2 E$ and ultimately E .

The moduli space of such instantons is then the total space of

$$\frac{SL(3, \mathbb{H})}{SU(3)} \longrightarrow \frac{SL(3, \mathbb{H})}{Sp(3)}, \text{ whose } T_o \cong \Lambda_0^2(\underline{\mathbb{C}}^6)$$

- In parallel to the theory of manifolds with reduced holonomy, there is a unified theory of instantons (which arguably preceded it in the exceptional cases).
- The quaternionic version of ADHM has many unanswered questions and still open problems [O]. Unlike for G_2 or $\text{Spin} 7$, there is no rich theory of submanifolds.
- The link between G_2 and the nearly-Kähler case is striking, especially since it is unknown if there are compact NK manifolds other than the four usual suspects (S^6 , $\mathbb{C}P^3$, \mathbb{F}^3 , $S^3 \times S^3$).
- More applications are needed of the differential complexes, their cohomology and index theory. A big problem is to determine h^2 .
- The exceptional cases (NK, G_2 , $\text{Spin} 7$) suffer from being non-holomorphic theories, with no twistor spaces to retire to. But this did not stop algebraic geometry being used in the construction of new compact manifolds with holonomy G_2 [CHNP].